# Velocity of Sound in Nitrogen and Argon at High Pressures 

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#### Abstract

Measurements of the velocity of sound in $\mathrm{N}_{2}$ and Ar are reported at the ice point and $294.26^{\circ} \mathrm{K}\left(70^{\circ} \mathrm{F}\right)$ and at pressures between 1 and 70 atm . The ratio of the specific heats as a function of pressure and temperature was calculated from the experimental results.


## INTRODUCTION

PRECISE measurements of the sound velocity in real gases at different temperatures and pressures offer a useful tool for studying the equation of state. Recently, Gyorog ${ }^{1}$ developed a generalized equation of state derived from experimental measurements of compressibility factor and enthalpy and internal energy deviations. To gain insight into the validity of this equation of state for $\mathrm{N}_{2}$ and Ar , velocity measurements were made in these gases at the ice point and

Table I. Sound velocity and ratio of specific heats in nitrogen.

| Tempera- <br> ture <br> $\left({ }^{\circ} \mathrm{K}\right)$ | Pressure <br> $(\mathrm{atm})$ | $V_{0}$ (exptl.) <br> $(\mathrm{m} / \mathrm{sec})$ | $V_{0}$ (Gyorog) <br> $(\mathrm{m} / \mathrm{sec})$ | $\gamma$ |
| :---: | ---: | :---: | :---: | :---: |
| 273.15 | 1.00 | 337.04 | 336.95 | 1.4029 |
|  | 10.00 | 338.17 | 338.01 | 1.4229 |
|  | 30.00 | 341.37 | 341.25 | 1.4660 |
|  | 50.00 | 345.88 | 345.76 | 1.5089 |
|  | 70.00 | 351.55 | 351.58 | 1.5483 |
| 294.26 | 1.00 | 349.72 | 349.76 | 1.4014 |
|  | 5.00 | 350.44 | 350.37 | 1.4097 |
|  | 10.00 | 351.22 | 351.20 | 1.4185 |
|  | 30.00 | 355.11 | 355.12 | 1.4552 |
|  | 50.00 | 359.98 | 360.09 | 1.4900 |
|  | 70.00 | 366.06 | 366.10 | 1.5246 |

$294.26^{\circ} \mathrm{K}$ and at pressures from 1 to 70 atm . Also, knowledge of the molecular weight and the specific heat, along with the Gyorog equation, enabled velocity values to be calculated. The experimental data are compared with these calculated values, as well as with the data of van Itterbeek. ${ }^{2,3}$

## METHOD

The experimental method is based on the apparatus developed by the author and described in a previous paper. ${ }^{4}$ The $\mathrm{N}_{2}$ and Ar had purity of $99.95 \%$ or better.

[^0]Temperature stability was achieved by running the tests in an air-conditioned room, the temperature of which did not change from $294.26^{\circ} \mathrm{K}$ by more than $\pm 2 \mathrm{C}^{\circ}$ in a $24-\mathrm{h}$ period. The entire sound tube was immersed in a water bath to ensure that no drastic changes in temperature (not more than $\pm 0.1 \mathrm{C}^{\circ}$ ) occurred during the test period. The ice-point temperature was controlled by immersing the sound tube in a bath of pure ice and distilled water.

Pressures of 1 and 5 atm were measured on a $130-\mathrm{in}$. manometer, while calibrated Heise absolute-pressure gauges measured the 10 - to $70-\mathrm{atm}$ points.
Extensive initial tests ${ }^{4}$ showed that the apparatus was compatible with the Helmholtz-Kirchoff (H-K) equation. Thus the $\mathrm{H}-\mathrm{K}$ constant for the tube ( $\beta_{T}$ ) was equal, within experimental accuracy, to the theoretical value ( $\beta$ ) and, therefore, could be calculated. The velocity of sound in the tube was then measured at one or more frequencies and corrected with the calculated $\beta$ to yield the Laplacian free-gas velocity $V_{0}$.
The sound velocities were also calculated with the aid of Eq. (1), using the generalized virial coefficients of Gyorog ${ }^{1}$ (or the generalized virial coefficients derived from the Lennard-Jones potential $)^{5,6}$ :

$$
\begin{equation*}
V_{0}=\left\{(R T / M)\left[\phi+(\Psi / x)\left(\gamma^{0}-1\right)\right]\right\}^{\frac{1}{2}}, \tag{1}
\end{equation*}
$$

where

$$
\begin{aligned}
& \phi=1+2 B^{*} \rho^{*}+3 C^{*} \rho^{* 2}+4 D^{*} \rho^{* 3}, \\
& \Psi=\left[1+\left(B^{*}+T^{*} \frac{d B^{*}}{d T^{*}}\right) \rho^{*}+\left(C^{*}+T^{*} \frac{d C^{*}}{d T^{*}}\right) \rho^{* 2}\right. \\
& \left.+\left(D^{*}+T^{*} \frac{d D^{*}}{d T^{*}}\right) \rho^{* 3}\right]^{2}, \\
& x=1-\left(\gamma^{0}-1\right)\left[\left(2 T^{*} \frac{d B^{*}}{d T^{*}}+T^{* 2} \frac{d^{2} B^{*}}{d T^{* 2}}\right) \rho^{*}\right. \\
& \left.+\left(T^{*} \frac{d C^{*}}{d T^{*}}+\frac{1}{2} T^{* 2} \frac{d^{2} C^{*}}{d T^{* 2}}\right) \rho^{* 2}+\left(\frac{2}{3} T^{*} \frac{d D^{*}}{d T^{*}}+\frac{1}{3} T^{* 2} \frac{d^{2} D^{*}}{d T^{* 2}}\right) \rho^{* 3}\right],
\end{aligned}
$$

$\gamma^{0}$ is the ratio of heat capacities for the ideal gas state;

[^1]Table II. Sound velocity and ratio of specific heats in argon.

| $\begin{aligned} & \text { Tempera- } \\ & \text { ture } \\ & \left({ }^{\circ} \mathrm{K}\right) \end{aligned}$ | Pressure (atm) | ) |  | $\gamma$ |
| :---: | :---: | :---: | :---: | :---: |
|  |  | $V_{0}($ exptl. $)$ $(\mathrm{m} / \mathrm{sec})$ | $\begin{gathered} V_{0}(\text { Gyorog }) \\ (\mathrm{m} / \mathrm{sec}) \end{gathered}$ |  |
| 273.15 | 1.00 | 308.14 | 308.06 | 1.6715 |
|  | 10.00 | 308.42 | 308.48 | 1.7030 |
|  | 30.00 | 310.02 | 309.95 | 1.7822 |
|  | 50.00 | 312.29 | 312.23 | 1.8647 |
|  | 70.00 | 315.33 | 315.42 | 1.9485 |
| 294.26 | 1.00 | 319.90 | 319.88 | 1.6714 |
|  | 10.00 | 320.49 | 320.53 | 1.6978 |
|  | 30.00 | 322.64 | 322.67 | 1.7628 |
|  | 50.00 | 325.65 | 325.48 | 1.8325 |
|  | 70.00 | 328.95 | 329.04 | 1.8985 |

$R, T, M$ are the gas constant, absolute temperature, and molecular weight, respectively; $B^{*}, C^{*}, D^{*}$ are the generalized virial coefficients (Refs. 1 and 5); and $\rho^{*}, T^{*}$ are the reduced density and reduced temperature. Equation (1) is the form of the generalized virial equation given by Hovi (but the fourth virial coefficient and its derivatives were added to Hovi's equation by the author).

The ratio of the heat capacities $\gamma$ as a function of pressure and temperatures was calculated from

$$
\begin{equation*}
\gamma=\frac{V_{0}^{2}}{(R / M) T\left(1+2 B^{*} \rho^{*}+3 C^{*} \rho^{* 2}+4 D^{*} \rho^{* 3}\right)} \tag{2}
\end{equation*}
$$



Fig. 1. Sound velocity in nitrogen as a function of pressure. $(X$, present data; ——, Gyorog's equation; $O$, van Itterbeek's data ${ }^{2}$.)

[^2]
## RESULTS AND CONCLUSIONS

The free-gas sonic velocities and the specific-heat ratios are listed in Tables I and II. The data of this paper, the calculated values from Gyorog's equation, and van Itterbeek's data are compared in Figs. 1 and 2. Note that the agreement between the present data $(X)$ and the values calculated by Gyorog's equation (solid line) is excellent in all cases. However, van Itterbeek's data for nitrogen at the ice point are lower than those of the present paper. The deviation increases from about $0.5 \mathrm{~m} / \mathrm{sec}$ at low pressures to about $2.3 \mathrm{~m} / \mathrm{sec}$ at high pressures.

Also presented in Fig. 1 is a comparison between velocity values for $\mathrm{N}_{2}$ calculated at $298.72^{\circ} \mathrm{K}$ by Gyorog's equation and those of van Itterbeek ${ }^{2}$ at the


Fig. 2. Sound velocity in argon as a function of pressure. ( $X$, present data; -_, Gyorog's equation; $O$, van Itterbeek's data ${ }^{3}$.)
same temperature. Here it is seen that the van Itterbeek data agree with the predicted values except for his highest test pressure.

In Fig. 2, the data of this paper for argon at the ice point are in very good agreement with those of van Itterbeek. ${ }^{3}$ But comparison between the calculated values for argon at $299.75^{\circ} \mathrm{K}$ and those of van Itterbeek at the same temperature shows some deviation from the predicted curve in the pressure range of 20 to 50 atm . The author's data at $294.26^{\circ} \mathrm{K}$, however, agree with the calculated values.

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